# Recovery of distributed quantum information from a quantum erasure 

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Physical Nano-Memories Signal and Information Processing Lab


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## Outline

- Double slit experiment
- Superdense coding
- Teleportation
- A curious example of measurement
- The problem of node failure from a network and recovery mechanism


## Double slit experiment



## Introduction

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- Physics allows us to manipulate information in a whole new setting.
- Encryption possible by making use of the postulates of quantum mechanics.


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- Measurement The act of measurement changes the state. This is unlike the classical world where the act of measuring the resistance of a resistor does not change the resistance.


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- Here we say that we have performed measurement in the $\{|0\rangle,|1\rangle\}$ basis


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- The state for multiple qubits resides in the space obtained by tensor product of the individual spaces:

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\mathcal{H} & =\mathcal{H}_{1} \otimes \mathcal{H}_{2} \otimes \cdots \mathcal{H}_{n} \\
|\psi\rangle & =\left|\psi_{1}\right\rangle \otimes\left|\psi_{2}\right\rangle \otimes \cdots\left|\psi_{n}\right\rangle
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- This phenomenon is called as entanglement


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- A preshared Bell pair lets A communicate two classical bits.


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- If $A$ measures its qubit in the $\{|+\rangle,| \rangle\}$ basis, then $B$ and $C$ end up in either $\frac{|00\rangle+|11\rangle}{\sqrt{2}}$ or $\frac{|00\rangle-|11\rangle}{\sqrt{2}}$


## The problem of node failure and recovery mechanism

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- We use the following setting for describing quantum networks via graphs.
- Nodes represented by vertices from the set $V$ account for qubits, one at every node, while edges represent the quantum mechanical interaction that enables entangling of qubits.



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Suppose $|\psi\rangle$ is a pure state of a composite system $A B$ with Hilbert space $\mathcal{H}_{A} \otimes \mathcal{H}_{B}$. Then there exist orthonormal states $|i\rangle_{A}$ for system $A$ and orthonormal states $|i\rangle_{B}$ for system $B$ such that

$$
|\psi\rangle=\sum_{i} \lambda_{i}\left|i_{A}\right\rangle\left|i_{B}\right\rangle
$$

where $\lambda_{i}$ are non-negative real numbers satisfying $\sum_{i} \lambda_{i}^{2}=1$ known as Schmidt co-efficients.

## Node failure



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- Suppose the purification procedure uses the basis states $\left|\theta_{0}\right\rangle,\left|\theta_{1}\right\rangle \in \mathcal{H}$ to combine the new qubit and the mixed state $\rho^{\prime}$, then we get the following purified state in $\mathcal{H}^{\otimes N}$ :


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$$
\begin{align*}
& \left|\psi^{\prime}\right\rangle=\lambda_{0}\left|\theta_{0}\right\rangle\left|\eta_{0}\right\rangle+\lambda_{1}\left|\theta_{1}\right\rangle\left|\eta_{1}\right\rangle,  \tag{1}\\
& \quad \text { or } \\
& \left|\psi^{\prime}\right\rangle=\lambda_{0}\left|\theta_{1}\right\rangle\left|\eta_{0}\right\rangle+\lambda_{1}\left|\theta_{0}\right\rangle\left|\eta_{1}\right\rangle . \tag{2}
\end{align*}
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- We are investing ways of detecting and identifying the node that goes through a failure.


## References


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## Questions and suggestions

Thanks for your time ${ }^{-}$

