

# Lecture-10: Many Random Variables

## 1 Qualitative view of Random Variables

Physical systems of particles can be modelled by assigning a random variable to state of each particles. Here we assume the particles have no interaction or interaction is minimal. Thus for  $N$  particles we have an  $N$  length vector  $x = (x_1, x_2, \dots, x_N)$  with probability

$$P_N(x) = P_N(X_1, \dots, x_N). \quad (1)$$

An instance of this would be Boltzmann distribution for a physical system with  $N$  degrees of freedom. Shannon Entropy of this distribution is

$$H_N = -\mathbb{E} \log P_N(x) \quad (2)$$

Typically  $H_N$  grows linearly with  $N$  for large  $N$ . So entropy per variable  $h_n = H_N/N$  has a finite limit; so we define a quantity  $h$ .

**Definition 1.1.**  $h : \lim_{N \rightarrow \infty} h_N$

We will be interested in defining the quantity

$$r(x) = 1/N \log [1/P_N(x)] \quad (3)$$

to characterize any particular realization  $\underline{x}$ . We would like to know how close is  $r(x)$  to  $h_N$ . In most cases  $r(x)$  is peaked around  $r = h_N$ . Often probability distribution of  $r(x)$  behave exponentially with

$$P\{r(x) \approx \rho\} = e^{-NI(\rho)} \quad (4)$$

where  $I(\rho)$  has minimum at  $\rho = h$ . and  $I(h) = 0$  From this observation and (3) we have

$$P_N(x) = e^{-Nh} \quad (5)$$

Total probability of realization  $x$  with  $r(x) \approx h$  is 1. This will imply that number of configurations with major contribution to probability is  $e^{Nh}$ . This in general is a small number when compared to total number of configurations  $\mathcal{X}^N$ .

$$\text{Number of typical configurations} = e^{Nh} \ll \mathcal{X}^N \quad (6)$$

This has a lot of implications. For example if we want to estimate an observable  $\mathcal{O}$  we can sample configuration and it will give good estimates because major part of total probability is distributed among typical configurations. But it takes order of  $e^{N(\log|\mathcal{X}|-h)}$  time to generate a sample. Monte Carlo methods provide better solution for this problem

## 2 Large deviations for independent variables

**Theorem 2.1.** (Sanov) Let  $s_1, \dots, s_N \in \mathcal{X}$  be  $N$  i.i.d random variables drawn from the probability distribution  $p(x)$ , and let  $K \subset \mathcal{M}(\mathcal{X})$  be a compact set of probability distributions over  $\mathcal{X}$ . If  $q$  is the type of  $(s_1, \dots, s_N)$  then

$$\mathbb{P}[q \in K] = e^{-ND(q^* \| p)} \quad (7)$$

where  $q^* = \arg \min_{q \in K} D(q \| p)$ , and  $D(q \| p)$  is the Kullback-Leibler divergence.

*Proof.* Now for a particular  $q'$

$$\begin{aligned}
(P(q_s = q')) &= \prod_{x \in \mathcal{X}} P(Nq'(x) \text{ times } x \text{ appears in } s) \\
&= \binom{N}{Nq'(x_1), Nq'(x_2), \dots, Nq'(x_{|\mathcal{X}|})} \prod_{x \in \mathcal{X}} p(x)^{Nq'(x)} \\
&= \binom{N!}{Nq'(x_1)!, Nq'(x_2)!, \dots, Nq'(x_{|\mathcal{X}|})!} \prod_{x \in \mathcal{X}} p(x)^{Nq'(x)}
\end{aligned} \tag{8}$$

Using Stirlings equation

$$\log_2 n! = n \log_2 n \tag{9}$$

$$n! = 2^{n \log_2 n} \tag{10}$$

$$\frac{N!}{(Nq'(x_1))!(Nq'(x_2))! \dots (Nq'(x_{|\mathcal{X}|}))!} = 2^{N \log N - Nq_1 \log Nq_1 - Nq_2 \log Nq_2 \dots Nq_{|\mathcal{X}|} \log Nq_{|\mathcal{X}|}} \tag{11}$$

$$\prod_{x \in \mathcal{X}} p(x)^{Nq'(x)} = \prod_{x \in \mathcal{X}} 2^{Nq'(x) \log p(x)} \tag{12}$$

We have KL Divergence D as

$$D(q \| p) = \sum_x q(x) \log \frac{q(x)}{p(x)} \tag{13}$$

From 8, 10, 11, 12 we have

$$P(q_s = q') = \prod_{x \in \mathcal{X}} 2^{-ND(q \| p)} \tag{14}$$

By approximating with leading term in exponential we have

$$P(q_s = q') = e^{-ND(q \| p)} \tag{15}$$

□

**Example 2.2.** A simple model of a column of the atmosphere is obtained by considering  $N$  particles in the earth's gravitational field. The state of particle  $i \in \{1, 2, \dots, N\}$  is given by a single coordinate  $z_i \geq 0$  which measures its height with respect to level. To simplify, we make assumption  $z_i$  is integer.

$$E = \sum_{i=1}^N Nz_i \tag{16}$$

Consider a configuration :  $z_1, \dots, z_N$ . Its type can be interpreted as:

$$\rho(z) = 1/N \sum_{i=1}^N \mathbb{1}_{\{z=z_i\}} \tag{17}$$

$$\rho_{eq}(z) = \langle \rho(z) \rangle = (1 - e^{-\beta}) e^{-\beta z} \tag{18}$$

Using Boltzmann probability distribution; expected density profile can be computed as follows.

$$\text{Partition function } Z(\beta) = \sum_{x \in \mathcal{X}} e^{-E(x)} = \sum_{z=0}^{\infty} e^{-\beta z} = \frac{1}{1 - e^{-\beta}}. \tag{19}$$

So probability distribution follows

$$\mu_\beta(x) = \frac{1}{Z(\beta)} e^{-\beta E(x)} = (1 - e^{-\beta}) e^{-\beta z} \tag{20}$$

We compute the probability of getting a general exponential density profile with parameter  $\lambda$ :

$$\rho_\lambda(z) = (1 - e^{-\lambda})e^{-\lambda z} \quad (21)$$

For that we find the KL divergence between  $\rho_\lambda$  and  $\rho_{eq}$

$$D(\rho_\lambda \parallel \rho_{eq}) = \sum_z \rho_\lambda(z) \log \frac{\rho_\lambda(z)}{\rho_{eq}(z)} = \log \frac{1 - e^{-\lambda}}{1 - e^{-\beta}} + \frac{\beta - \lambda}{e^\lambda - 1} \quad (22)$$

We now plot KL divergence as a function of  $\lambda$  in figure 1.

$$I_\beta(\lambda) = D(\rho_\lambda \parallel \rho_{eq}) \quad (23)$$

It can be noted from the figure 1 that small values of  $\lambda$  are very rare.

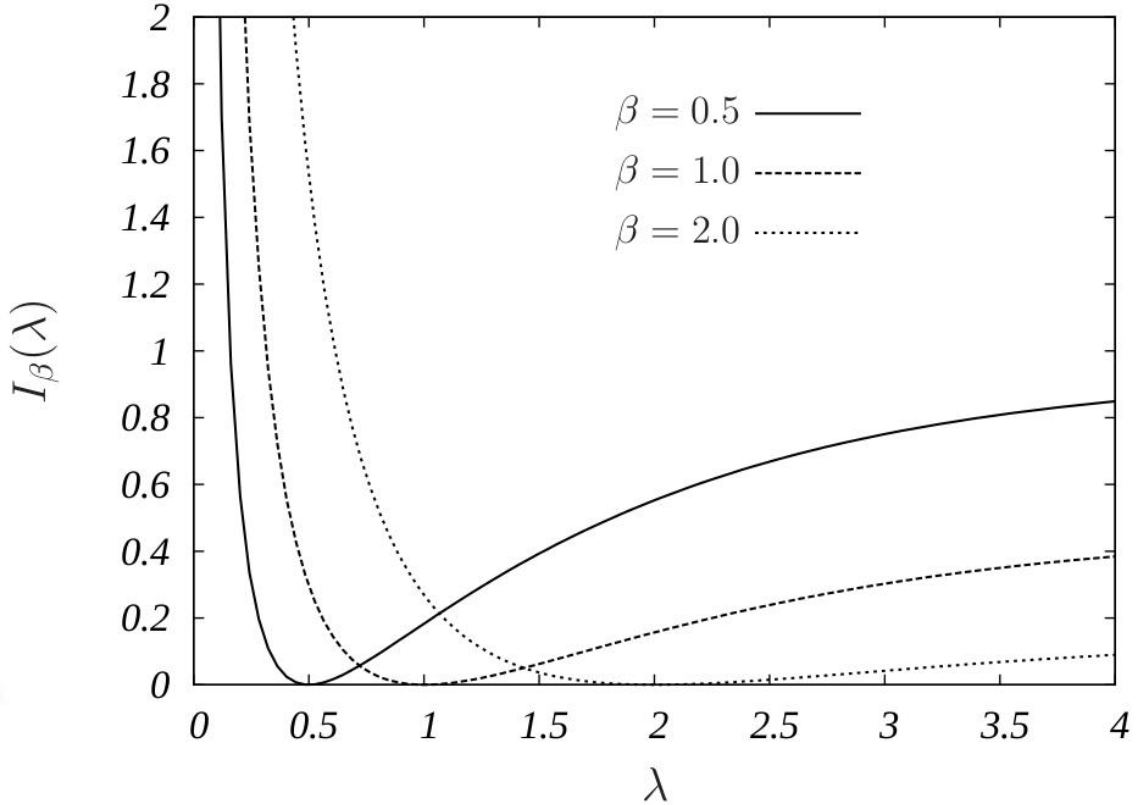


Figure 1: Equilibrium density profile

### 3 How typical is an empirical average?

The empirical average of a measurement is given by:

$$\bar{f} = \frac{1}{N} \sum_{i=1}^N f(s_i) \quad (24)$$

**Corollary 3.1.** Let  $s_1, \dots, s_N$  be  $N$  i.i.d random variables drawn from a probability distribution  $p(\cdot)$ . Let  $f : \mathcal{X} \rightarrow \mathbb{R}$  be a real-valued function and let  $\bar{f}$  be its empirical average. If  $A \subset \mathbb{R}$  is a closed interval of the real axis, then

$$\mathbb{P}[\bar{f} \in A] = e^{[-NI(A)]}, \quad (25)$$

where

$$I(A) = \min_q [D(q||p) | \sum_{x \in \mathcal{X}} q(x)f(x) \in A] \quad (26)$$

*Proof.* We note that  $\bar{f}$  is related to the type of sequence  $x_1, x_2, \dots, x_n$  as  $\bar{f} = \sum_x q(x)f(x)$ . Keeping in mind Sanov theorem, we define a set as follows

$$K = \left\{ q \in \mathcal{M}(\mathcal{X}) \mid \sum_{x \in \mathcal{X}} q(x)f(x) \in A \right\} \quad (27)$$

Then the result follows directly from Sanov by comparing  $K$  with set in Sanov and  $I(A)$  with  $D(q^*||p)$ .  $\square$

**Example 3.2.** Let  $s_1, \dots, s_N$  be  $N$  i.i.d random variables drawn from a probability distribution  $p(\cdot)$  with bounded support. Show that, to leading exponential order,

$$\mathbb{P}\{s_1 + \dots + s_N \leq 0\} = \left\{ \inf_{z \geq 0} \mathbb{E} e^{-zs_1} \right\}^N \quad (28)$$

**Example 3.3.** We look at  $N$  particles in a gravitational field, and consider the average height of the particles

$$\bar{z} = \frac{1}{N} \sum_{i=1}^N z_i. \quad (29)$$

Expected value of this quantity is

$$\mathbb{E}\bar{z} = z_{eq} = (e^\beta - 1)^{-1}. \quad (30)$$

The probability of a fluctuation in  $\bar{z}$  is easily computed using the above corollary. For  $z > z_{eq}$  we obtain

$$\mathbb{P}(z > \bar{z}) = e^{-NI(z)}, \quad (31)$$

where

$$I(z) = (1+z) \log \left( \frac{1+z_{eq}}{1+z} \right) + z \log \left( \frac{z}{z_{eq}} \right) \quad (32)$$